

The meaning of H

Suppose Ω is a region in $\mathbb{R}^2$ with smooth boundary $\partial\Omega$ and

$$\vec{X}(x,y) = (x, y, z(x,y))$$

We have previously computed

$$\vec{X}_x = (1, 0, z_x)$$

$$\vec{X}_y = (0, 1, z_y)$$

so

$$E = 1 + z_x^2, \quad F = z_x z_y, \quad G = 1 + z_y^2$$

Further,

$$\vec{X}_x \times \vec{X}_y = (-z_x, -z_y, 1)$$

so

$$\vec{n} = \frac{(-z_x, -z_y, 1)}{\sqrt{1 + z_x^2 + z_y^2}}$$

This means

$$l = \langle \vec{x}_{xx}, \vec{n} \rangle = \frac{z_{xx}}{\sqrt{1 + z_x^2 + z_y^2}}$$

$$m = \langle \vec{x}_{xy}, \vec{n} \rangle = \frac{z_{xy}}{\sqrt{1 + z_x^2 + z_y^2}}$$

$$n = \langle \vec{x}_{yy}, \vec{n} \rangle = \frac{z_{yy}}{\sqrt{1 + z_x^2 + z_y^2}}$$

We want to compute the trace of the shape operator. Since

$$S_p = \begin{bmatrix} E & F \\ F & G \end{bmatrix}^{-1} \begin{bmatrix} l & m \\ m & n \end{bmatrix}$$

we can compute

$$\begin{aligned} S_p &= \frac{1}{EG - F^2} \begin{bmatrix} G & -F \\ -F & E \end{bmatrix} \begin{bmatrix} l & m \\ m & n \end{bmatrix} \\ &= \frac{1}{EG - F^2} \begin{bmatrix} Gl - Fm & \sim \\ \sim & -Fm + En \end{bmatrix} \end{aligned}$$

so

$$\text{tr } \mathcal{J}_p = \frac{Gl - 2Fm + En}{EG - F^2}$$

and

$$H = \frac{1}{2} \text{tr } \mathcal{J}_p = \frac{Gl - 2Fm + En}{2(EG - F^2)}$$

Now

$$\begin{aligned} EG - F^2 &= (1 + z_x^2)(1 + z_y^2) - z_x^2 z_y^2 \\ &= 1 + z_x^2 + z_y^2 \end{aligned}$$

In the numerator, we get

$$Gl - 2Fm + En =$$

$$\frac{(1 + z_y^2)z_{xx} - 2z_x z_y z_{xy} + (1 + z_x^2)z_{yy}}{\sqrt{1 + z_x^2 + z_y^2}}$$

Theorem. For a graph surface $(x, y, z(x, y))$

$$H = \frac{(1 + z_y^2) z_{xx} - 2 z_x z_y z_{xy} + (1 + z_x^2) z_{yy}}{2(1 + z_x^2 + z_y^2)^{3/2}}$$

We now want to interpret this!

$$\begin{aligned} \text{Area}(z) &= \int_{\Omega} \sqrt{EG - F^2} \, dx \, dy \\ &= \int_{\Omega} \sqrt{1 + z_x^2 + z_y^2} \, dx \, dy \\ &= \int_{\Omega} \sqrt{1 + \|\nabla z\|^2} \, dx \, dy \end{aligned}$$

Now suppose we want to compute the directional derivative

$$D_v \text{Area}(z) = \frac{d}{dt} \text{Area}(z + tv)$$

for some variation $v: \Omega \rightarrow \mathbb{R}$ with $v \equiv 0$ on $\partial\Omega$.

We can write down

$$\left. \frac{d}{dt} \text{Area}(z+tv) \right|_{t=0} =$$

$$= \int_{\Omega} \left. \frac{d}{dt} \sqrt{1 + (z+tv)_x^2 + (z+tv)_y^2} \right|_{t=0} dx dy$$

Now

$$\begin{aligned} \left. \frac{d}{dt} (z+tv)_x^2 \right|_{t=0} &= \left. \frac{d}{dt} (z_x + tv_x)^2 \right|_{t=0} \\ &= 2(z_x + tv_x) \cdot v_x \Big|_{t=0} \\ &= 2z_x v_x \end{aligned}$$

Similarly,

$$\left. \frac{d}{dt} (z+tv)_y^2 \right|_{t=0} = 2z_y v_y.$$

Therefore

$$\left. \frac{d}{dt} \sqrt{1 + (z+tv)_x^2 + (z+tv)_y^2} \right|_{t=0} =$$

$$= \frac{1}{\cancel{2} \sqrt{1 + z_x^2 + z_y^2}} \cdot (\cancel{2} z_x v_x + \cancel{2} z_y v_y)$$

$$= \left\langle \frac{\nabla z}{\sqrt{1 + \|\nabla z\|^2}}, \nabla v \right\rangle$$

Now in vector calculus, we learn to define the divergence of a vector field $\vec{F} : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ by

$$\operatorname{div} \vec{F} = \frac{\partial}{\partial x_1} F_1 + \frac{\partial}{\partial x_2} F_2$$

Thus, for any scalar function ϕ ,

$$\operatorname{div} \phi \vec{F} = \frac{\partial}{\partial x_1} \phi F_1 + \frac{\partial}{\partial x_2} \phi F_2$$

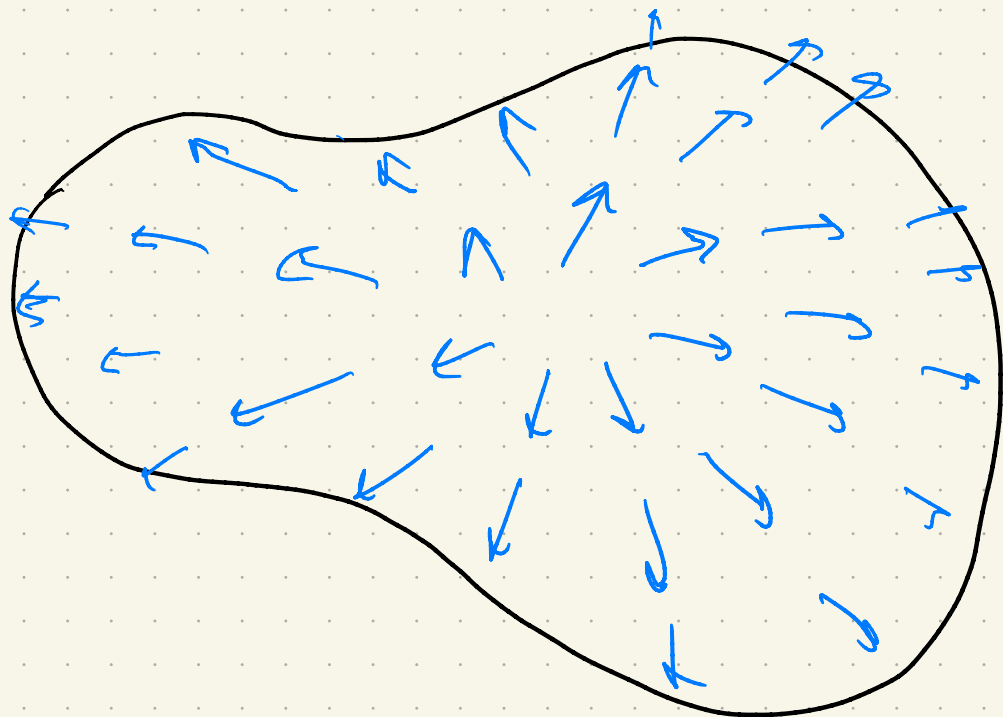
$$= \left(\frac{\partial}{\partial x_1} \varphi \right) F_1 + \varphi \left(\frac{\partial}{\partial x_1} F_1 \right) + \left(\frac{\partial}{\partial x_2} \varphi \right) F_2 + \varphi \left(\frac{\partial}{\partial x_2} F_2 \right)$$

$$= \langle \nabla \varphi, F \rangle + \varphi \operatorname{div} F.$$

We also learn the divergence theorem,

$$\int_{\Omega} \operatorname{div} F \, d\text{Area} = \int_{\partial\Omega} \langle F, \vec{n} \rangle \, ds$$

where $\vec{n}$ is the outward pointing unit normal to $\partial\Omega$.



Putting these together,

$$\int_{\Omega} \left\langle \frac{\nabla z}{\sqrt{1 + \|\nabla z\|^2}}, \nabla v \right\rangle d\text{Area} =$$

$$= \int_{\Omega} \operatorname{div} \left(v \frac{\nabla z}{\sqrt{1 + \|\nabla z\|^2}} \right) - v \operatorname{div} \left(\frac{\nabla z}{\sqrt{1 + \|\nabla z\|^2}} \right) d\text{Area}$$

$$= \int_{\partial\Omega} \left\langle v \frac{\nabla z}{\sqrt{1 + \|\nabla z\|^2}}, \vec{n} \right\rangle ds - \int_{\Omega} v \operatorname{div} \left(\frac{\nabla z}{\sqrt{1 + \|\nabla z\|^2}} \right) d\text{Area}$$

0, because $v=0$ on $\partial\Omega$

So $D_v \text{Area}(z) = 0$ for all variations v
if and only if $\operatorname{div} \left(\frac{\nabla z}{\sqrt{1 + \|\nabla z\|^2}} \right) \equiv 0$.

So now let's compute

$$\operatorname{div} \left(\frac{1}{\sqrt{1+z_x^2+z_y^2}} \nabla z \right) = \left\langle \nabla \left(\frac{1}{\sqrt{1+z_x^2+z_y^2}} \right), \nabla z \right\rangle + \frac{z_{xx} + z_{yy}}{\sqrt{1+z_x^2+z_y^2}}$$

Now

$$\nabla \left(\frac{1}{\sqrt{1+z_x^2+z_y^2}} \right) = -\frac{1}{\cancel{2} (1+z_x^2+z_y^2)^{3/2}} \begin{bmatrix} \cancel{2} z_x z_{xx} + \cancel{2} z_y z_{yx} \\ \cancel{2} z_x z_{xy} + \cancel{2} z_y z_{yy} \end{bmatrix}$$

So

$$\begin{aligned} & \left\langle \nabla \frac{1}{\sqrt{1+z_x^2+z_y^2}}, \nabla z \right\rangle + \frac{z_{xx} + z_{yy}}{\sqrt{1+z_x^2+z_y^2}} \\ &= \frac{-z_x^2 z_{xx} - z_x z_y z_{yx} - z_x z_y z_{xy} - z_y^2 z_{yy}}{(1+z_x^2+z_y^2)^{3/2}} \\ &+ \frac{(1+z_x^2+z_y^2)(z_{xx} + z_{yy})}{(1+z_x^2+z_y^2)^{3/2}} \end{aligned}$$

$$= \frac{(1+z_y^2)z_{xx} - 2z_x z_y z_{xy} + (1+z_x^2)z_{yy}}{(1+z_x^2+z_y^2)^{3/2}}$$

$$+ \frac{z_x^2 z_{xx} - \cancel{z_x^2 z_{xx}} + z_y^2 z_{yy} - \cancel{z_y^2 z_{yy}}}{(1+z_x^2+z_y^2)^{3/2}}$$

which proves a striking theorem!

Theorem. $2H$ is the first variation of area.

Surfaces with $H \equiv 0$ are area-critical
we call them minimal surfaces.

Example. $(x, y, \log(\frac{\cos x}{\cos y}))$.

We compute

$$Z(x, y) = \log(\cos x) - \log(\cos y)$$

so

$$Z_x = \frac{1}{\cos x} \cdot (-\sin x) = -\tan x$$

$$Z_y = -\frac{1}{\cos y} (-\sin y) = \tan y$$

and

$$Z_{xx} = -\sec^2 x$$

$$Z_{xy} = 0$$

$$Z_{yy} = -\sec^2 y$$

so

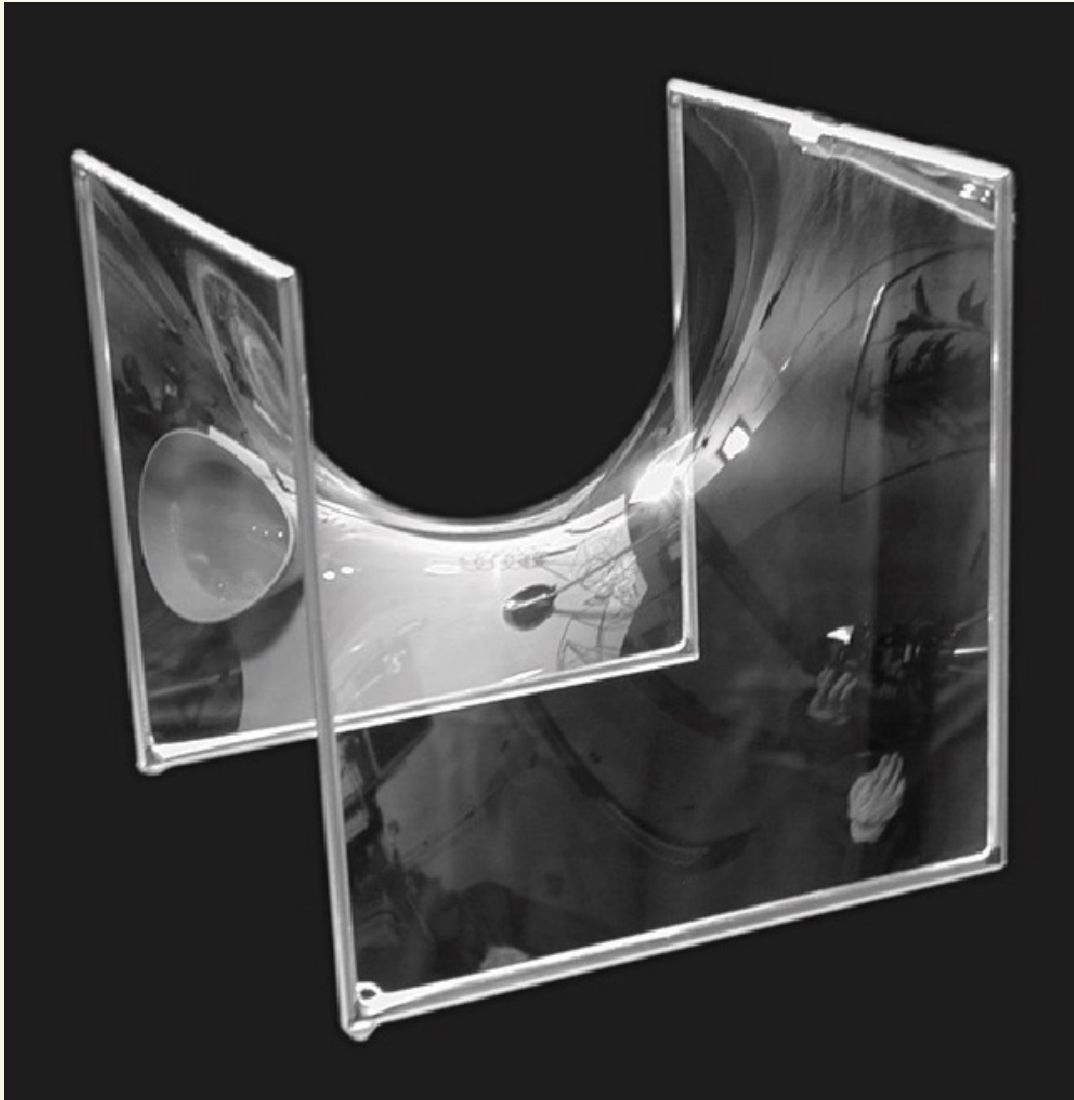
$$H = \frac{(1 + z_y^2) z_{xx} - 2 z_x z_y z_{xy} + (1 + z_x^2) z_{yy}}{2(1 + z_x^2 + z_y^2)^{3/2}}$$
$$= \frac{-(1 + \tan^2 y) \sec^2 x + (1 + \tan^2 x) \sec^2 y}{2(1 + \tan^2 x + \tan^2 y)^{3/2}}$$

But $1 + \tan^2 \theta = \sec^2 \theta$ for all θ , so

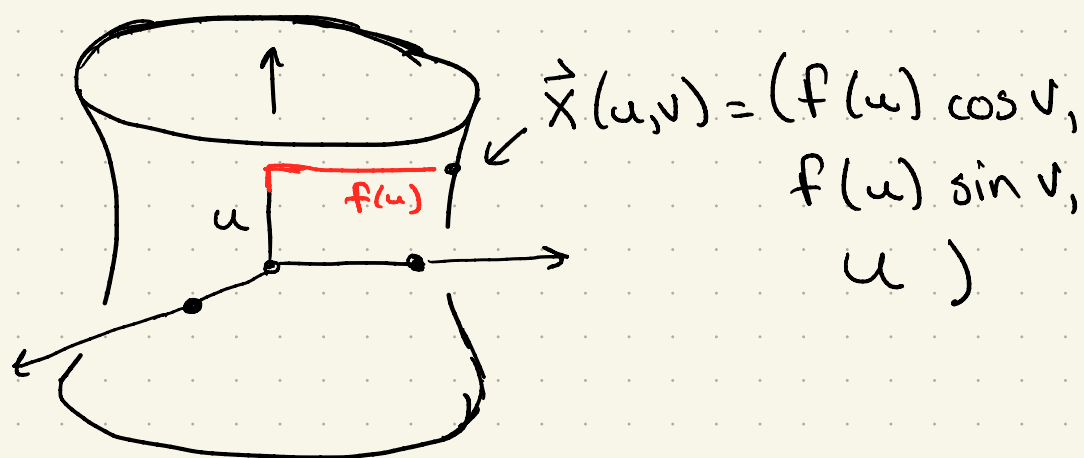
$$= \frac{-\sec^2 y \sec^2 x + \sec^2 x \sec^2 y}{2(1 + \tan^2 x + \tan^2 y)^{3/2}}$$
$$= 0$$

and this is well-defined as long as $x, y \in (-\pi/2, \pi/2)$ so $\cos x, \cos y \neq 0$.

This is called Scherk's surface.



Example. In previous notes, we considered the surface of revolution



and showed that

$$H = \frac{(1 + (f')^2) - f f''}{2f(1 + (f')^2)^{3/2}} \text{ sign}(f)$$

This means that $H=0 \iff$

$$(1 + (f')^2) - f f'' = 0$$

$$f'' = \frac{1 + (f')^2}{f}$$

This is a 2nd order ODE for f .

We can solve it by guessing and checking.

$$f(u) = c \cosh \frac{u}{c}$$

$$f'(u) = \sinh \frac{u}{c}$$

$$f''(u) = \frac{1}{c} \cosh \frac{u}{c}$$

So

$$\frac{1 + f'(u)^2}{f(u)} = \frac{1 + \sinh^2 \frac{u}{c}}{c \cosh \frac{u}{c}} = \frac{\cosh^2 \frac{u}{c}}{c \cosh \frac{u}{c}} = f''(u)$$

(In fact, I think this is the only real solution, as $f(u) = iu$ also solves the ODE.)

The resulting surface is the catenoid!

